

# Very Special Relativity as a background field theory

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- 1605.04967 [hep-th], Phys. Rev. D 94 (2016) 045019

# Outline

1. Intro: Lorentz violation and Very Special Relativity
2. Intro: Strong Field QED.
3. Particle motion in VSR and in background fields
4. SF-QED  $\leftrightarrow$  SIM(2)-QED correspondence

# What is VSR?

- New physics from violation of spacetime symmetries?

Liberati, *Class. Quant. Grav.* 30 (2013) 133001

## VSR

1. Replace Lorentz group (6 dim)  $\rightarrow$  subgroup (2,3,4 dim.)
2. Keep translation invariance.
3. Build (Lorentz invariance violating) QFTs.

Cohen & Glashow, *PRL* 97 (2006) 021601

- Original motivation: allows for neutrino masses.

Cohen & Glashow, *hep-ph/0605036*

## SIM(2)

- Largest subgroup: 4-dimensional SIM(2).

$$J_3, \quad K_3, \quad T_1 = K^1 + J^2, \quad T_2 = K^2 - J^1$$


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- $c$  constant, same dispersion.  $\rightarrow p^2 = \text{constant}$ .
- SIM(2) + parity = Lorentz.  $\rightarrow$  Parity violation.  
B. Bucher et al, PRL 116 (2016) 112503
- Null hyperplanes invariant...  $\rightarrow$  Lightlike stability group
- Preferred lightlike direction  $\rightarrow$  From an æther?  
Gibbons et al, PRD 76 (2007) 081701  
Cheon et al, PLB 679 (2009) 73

Direction, but no velocity

# SIM(2) and mass terms

- SIM(2) permits **nonlocal** terms of the form  $\frac{n_\mu}{n \cdot \partial}$

Dirac equation. . .

$$\not{p} = m \quad \longrightarrow$$

. . . in SIM(2)

$$\not{p} - \frac{\delta m^2}{2n \cdot p} \not{n} = m$$

Mass-shell:

$$p^2 = m^2 \quad \longrightarrow$$

$$p^2 = m^2 + \delta m^2 .$$

- $\delta m \leftrightarrow$  **small** neutrino mass.

Cohen & Glashow PRL 97 (2006) 021601

# Particle motion in background fields

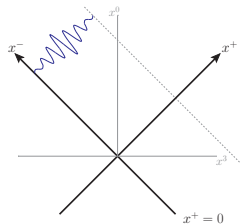
- Background plane wave field (basic laser model):

$$eA_\mu(x) =: a_\mu(n \cdot x) = ma_0(0, \cos(\omega n \cdot x), \sin(\omega n \cdot x), 0)$$

$$n_\mu = (1, 0, 0, 1)$$

- Lightlike wave vector  $n_\mu$ .
- Frequency  $\omega$ .
- Intensity (coupling)  $a_0$ .
- Integrable model

(Lorentz, KG, Dirac all solvable)



# Classical motion

- $\sigma$ , rates involve **averages** over many field oscillations.

Nikishov & Ritus 1963

- Example: classical momentum.

$$\pi_{\mu}(n.x) = p_{\mu} - a_{\mu}(n.x) + \frac{2p.a(n.x) - a^2(n.x)}{2n.p} n_{\mu}$$

**average**  $\downarrow$  (linear in  $a \rightarrow 0$ , quadratic  $\rightarrow 0$ )

$$q_{\mu} := \langle \pi_{\mu} \rangle = p_{\mu} + \frac{a_0^2 m^2}{2n.p} n_{\mu} .$$



# VSR connection



$$\pi_\mu \longrightarrow q_\mu = p_\mu + \frac{a_0^2 m^2}{2n \cdot p} n_\mu \quad \text{quasimomentum}$$

$$\pi^2 = m^2 \longrightarrow q^2 = m^2 + a_0^2 m^2 \quad \text{effective mass}$$

T.W.B. Kibble, Phys. Rev. 138 (1965) B740,

Harvey, Heinzl, Ilderton, Marklund PRL 109 (2012) 100402

- Frequency scale drops out
- **Averaged** motion in plane waves described by VSR?

$$a_0^2 m^2 \leftrightarrow \delta m^2$$

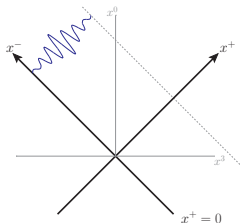


# QED in background fields

- QED in a background field / BSM photon background.

$$eA_\mu(x) =: a_\mu(n \cdot x) = ma_0(0, \cos(\omega n \cdot x), \sin(\omega n \cdot x), 0)$$
$$n_\mu = (1, 0, 0, 1)$$

- ★ Very high frequency scale  $\omega$ .
- ★ Cannot be probed by typical processes.
- ★ Effective physical observables: averages.



# Averaging

## Basic idea.

- Obtain an “effective” theory by:

1. **Killing** rapidly oscillating terms

$$\text{e.g. } a_\mu(n.x)$$

2. **Keeping** slowly varying terms

$$\text{e.g. } a^2(n.x) = -\delta m^2$$

- Classically: gives quasi-momentum/VSR momentum.
- Quantum: consistent, gauge invariant truncation of QED?

# Dirac equation

- Dirac equation in the background plane wave:

$$(i\cancel{\partial} - \cancel{a})\psi - m\psi = 0$$

- Volkov solutions:

$$\psi = \int dp \left( 1 + \frac{\cancel{a}(n \cdot x)}{2n \cdot p} \right) u_p \exp \left[ -ip \cdot x - i \int^{n \cdot x} 2p \cdot a - a^2 \right] b_p + \dots$$

- Average:**  $\psi \rightarrow \psi_{\text{av}}$ . The averaged field obeys:  $a_0^2 m^2 \leftrightarrow \delta m^2$

$$\left( i\cancel{\partial} - \frac{\delta m^2}{2i n \cdot \cancel{\partial}} \cancel{a} \right) \psi_{\text{av}} - m\psi_{\text{av}} = 0$$

- The Dirac equation in SIM(2) VSR!

# Averaging in QED

- Fermion propagator in a plane wave. [D.M. Volkov, Z. Phys. 94 \(1935\) 250](#)  
[H. Mitter, Acta Phys. Austr. Suppl. 14 \(1975\) 397](#)

$$S_{\text{Volk}}(x, y) = \int dp e^{-i\dots} \left( 1 + \frac{\not{n}\not{d}(n.x)}{2n.p} \right) \frac{\not{p} + m}{p^2 - m^2 + i\epsilon} \left( 1 + \frac{\not{d}\not{n}(n.y)}{2n.p} \right)$$

# Averaging in QED

- Fermion propagator in a plane wave. D.M. Volkov, Z. Phys. 94 (1935) 250  
H. Mitter, Acta Phys. Austr. Suppl. 14 (1975) 397

$$S_{\text{Volk}}(x, y) \xrightarrow{\text{average}} S_{\text{vsr}}(x-y) = \int \frac{d^4q}{(2\pi)^4} \frac{\not{q} + m - \frac{\delta m^2 \not{q}}{2n \cdot q}}{q^2 - m^2 - \delta m^2 + i\epsilon} e^{-iq \cdot (x-y)}$$

- Restores translation invariance, as in VSR. ✓

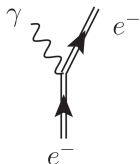
Cohen & Glashow PRL 97 (2006) 021601

- Poles shifted to VSR mass. ✓
- Spin- $\frac{1}{2}$  propagator of VSR. ✓

# The VSR vertex and the Ward Identity

☹ However ...

- Extra slowly-varying terms at **vertices** from products of fields.



$$\int \dots \left( 1 + \frac{\not{d}(n \cdot x) \not{n}}{2n \cdot p'} \right) \gamma^\mu \left( 1 + \frac{\not{n} \not{d}(n \cdot x)}{2n \cdot p} \right) G_{\mu\nu} \dots$$

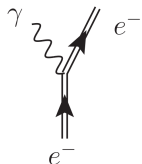
- Gives correction to the QED vertex:

$$\gamma^\mu \rightarrow \Gamma^\mu := \gamma^\mu + \frac{\delta m^2 n^\mu}{2n \cdot (p + k)n \cdot p}$$

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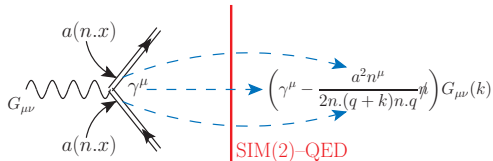
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☺ The three-point vertex of VSR. ✓

# The VSR vertex and the Ward Identity

- Provided **all** slowly varying terms are retained

Ilderton, PRD 94 (2016) 045019



1. VSR 3-point vertex recovered.
2. Ward Identity is preserved:

$$k \cdot \Gamma = 0 \quad \checkmark$$



# Higher-order vertices

☹ **However...** VSR has  $n$ -photon vertices for all  $n$ .

Dunn & Mehen, hep-ph/0610202

- Back to the propagator:

$$S_{\text{VolK}}(x, y) = \int \dots \left( 1 + \frac{\not{n}\not{a}(n.x)}{2n.p} \right) \dots \not{p} \dots \left( 1 + \frac{\not{a}(n.y)\not{n}}{2n.p} \right)$$

- $n.x \neq n.y \implies a(n.x)a(n.y)$  **rapidly oscillating**.

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- $n.x \neq n.y \implies a(n.x)a(n.y)$  **rapidly oscillating**.
- What happens if  $n.x = n.y$ ?

→ **Short distance** behaviour of the propagator.

# Higher-order vertices

- Singular term in Volkov propagator:

$$S_{\text{Volk}} \supset i \int \frac{d^4 p}{(2\pi)^4} \frac{\not{n}}{2n \cdot p} e^{-ip \cdot (x-y)} \propto \delta(n \cdot x - n \cdot y) \not{n}.$$

- 'Instantaneous lightfront propagator'.

Brodsky et al, Phys.Rept. 301 (1998), Bakker . . . A.I. et al, Nucl.Phys.Proc.Suppl 251 (2014) 165

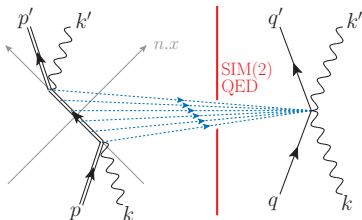
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- Generates precisely the higher-order VSR vertices. ✓

# VSR QED ↔ SFQED

- 
- VSR: an effective description of QED in a high frequency wave
- 
- What is the VSR 'æther'?
  - Preferred direction?
  - Nonlocal...
  - ...but trans. invariant.
  - Background EM wave.
  - Wave vector.
  - Interaction with the wave.
  - Restored by averaging.

# Conclusions

## Summary

- SIM(2) QED as an effective theory of SFQED.
- New connections between
  - ▶ BSM physics with/without Lorentz violation.
  - ▶ SFQED and Lorentz invariance violation

## Work in progress

- Other VSR groups?
- Other sectors of the standard model?
- Neutrino mass in this picture? (From e.g. loop effects.)
- Laser physics as “VSR in the lab”?

Heinzl and Ilderton, to appear